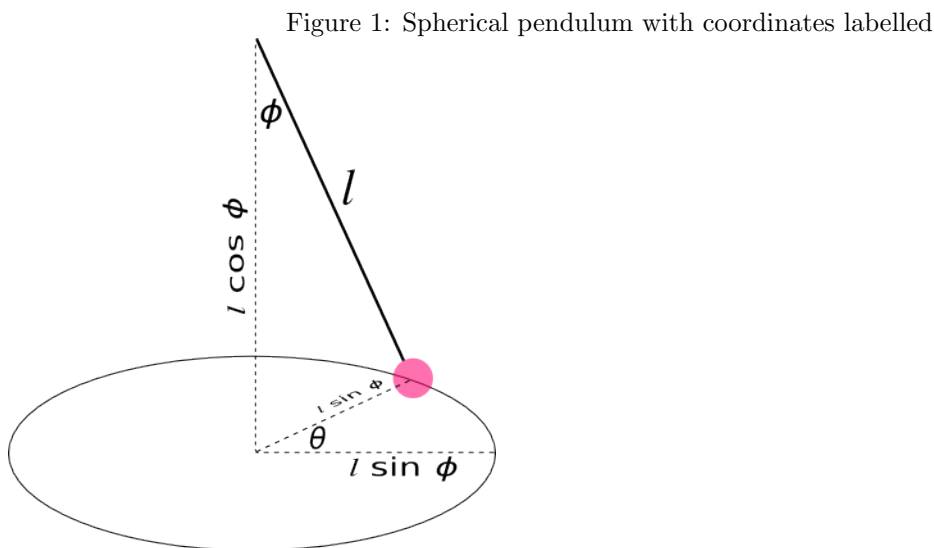


Goldstein, Poole, and Safko Problem 1.19

Josh Orndorff
admin@joshorndorff.com

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Establishing the coordinate system as it is labelled in the figure, we can begin by writing the potential and kinetic energy functions.

$$U = -mgl \cos \phi \quad (1)$$

$$T = \frac{1}{2}m \left(l^2 \dot{\theta}^2 \sin^2 \phi + l^2 \dot{\phi}^2 \right) \quad (2)$$

$$T = \frac{ml^2}{2} \left(\dot{\theta}^2 \sin^2 \phi + \dot{\phi}^2 \right) \quad (3)$$

That being established, we can write the Lagrangian as $L = T - U$.

$$L = \frac{ml^2}{2} \left(\dot{\theta}^2 \sin^2 \phi + \dot{\phi}^2 \right) + mgl \cos \phi \quad (4)$$

Now that we have the Lagrangian in all of its glory, we can compute the appropriate derivatives that we will need to find the equations of motion.

$$\begin{aligned} \frac{\partial L}{\partial \theta} &= 0 & \frac{\partial L}{\partial \phi} &= ml^2 \dot{\theta}^2 \sin \phi \cos \phi \\ \frac{\partial L}{\partial \dot{\theta}} &= ml^2 \dot{\theta} \sin^2 \phi & \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} &= ml^2 \dot{\phi} \\ \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} &= ml^2 \left(\ddot{\theta} \sin^2 \phi + 2\dot{\theta} \dot{\phi} \sin \phi \right) & \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} &= ml^2 \ddot{\phi} \end{aligned}$$

Now we can use the Euler Lagrange equation to determine the equations of motion. The ϕ equation becomes:

$$ml^2 \dot{\theta}^2 \sin \phi \cos \phi - mgl \sin \phi = ml^2 \ddot{\phi} \quad (5)$$

$$\ddot{\phi} - \dot{\theta}^2 \sin \phi \cos \phi + \frac{g}{l} \sin \phi = 0 \quad (6)$$

The Euler Lagrange equation for */theta* is:

$$ml^2 \left(\ddot{\theta} \sin^2 \phi + 2\dot{\theta}\dot{\phi} \sin \phi \right) = 0 \quad (7)$$

What this really tells us is that the quantity $\frac{\partial L}{\partial \dot{\theta}}$ is conserved.

These equations are not solvable analytically, but could be solved numerically.